

First JAM results

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(for the collaboration)



Data considered at this (first) stage

World data on polarized DIS

PROTON	SLAC E80/E130: G. Baum et al., Phys. Rev. Lett. 51 , 1135 (1983) EMC: J. Ashman et al., Nucl. Phys. B328 , 1 (1989) SMC: B. Adeva et al., Phys. Rev. D 58 , 112001 (1998) Phys. Rev. D 60 , 072004 (1999) Erratum-ibid. Phys. Rev. D 62 , 079902 (2000) COMPASS: M.G. Alekseev et al., Phys. Lett. B 690 , 466 (2010) SLAC E143: K. Abe et al., Phys. Rev. D 58 , 112003 (1998) SLAC E155: P.L. Anthony et al., Phys. Lett. B 493 , 19 (2000) Phys. Lett. B 458 , 529 (2000) SLAC E155x: P.L. Anthony et al., Phys. Lett. B 553 , 18 (2003) HERMES: A. Airapetian et al., Phys. Rev. D 75 , 012007 (2007) JLab Hall B (EG1b): Y. Prok et al., Phys. Lett. B 672 , 12 (2009) HERMES: A. Airapetian et al., Eur. Phys. J. C 72 , 1921 (2012)
DEUTERON	SMC: B. Adeva et al., Phys. Rev. D 58 , 112001 (1998) Phys. Rev. D 60 , 072004 (1999) Erratum-ibid. Phys. Rev. D 62 , 079902 (2000) COMPASS: V.Yu. Alexakhin et al., Phys. Lett. B 647 , 8 (2007) SLAC E143: K. Abe et al., Phys. Rev. D 58 , 112003 (1998) SLAC E155: P.L. Anthony et al., Phys. Lett. B 463 , 339 (1999) Phys. Lett. B 458 , 529 (2000) SLAC E155x: P.L. Anthony et al., Phys. Lett. B 553 , 18 (2003) HERMES: A. Airapetian et al., Phys. Rev. D 75 , 012007 (2007) JLab Hall B (EG1b): Y. Prok et al., Phys. Lett. B 672 , 12 (2009)
HELIUM-3	SLAC E142: P.L. Anthony et al., Phys. Rev. D 54 , 6620 (1996) SLAC E154: K. Abe et al., Phys. Rev. Lett. 79 , 26 (1997) Yu. Kolomensky, Ph.D. thesis, U. Massachusetts (1997), SLAC-Rep-503 HERMES: K. Ackerstaff et al., Phys. Lett. B 404 , 383 (1997) JLab Hall A (E99-117): X. Zhang et al., Phys. Rev. Lett. 92 , 012004 (2004) Phys. Rev. C 70 , 065207 (2004) JLab Hall A (E97-103): K. Kramer et al., Phys. Rev. Lett. 95 , 142002 (2005) K. Kramer, Ph.D. thesis, Coll. of William & Mary (2003) JLab Hall A (E01-012): P. Solvignon et al., Phys. Rev. Lett. 101 , 182502 (2008)

Mainly on *measured* asymmetries:

$$A_{\parallel} = D(A_1 + \eta A_2)$$

$$A_{\perp} = d(A_2 - \xi A_1)$$

D, d depend on

$$R = \frac{F_L}{(1 + \gamma^2)F_2 - F_L}$$

$$\gamma^2 = 4 \frac{M^2}{Q^2} x^2$$

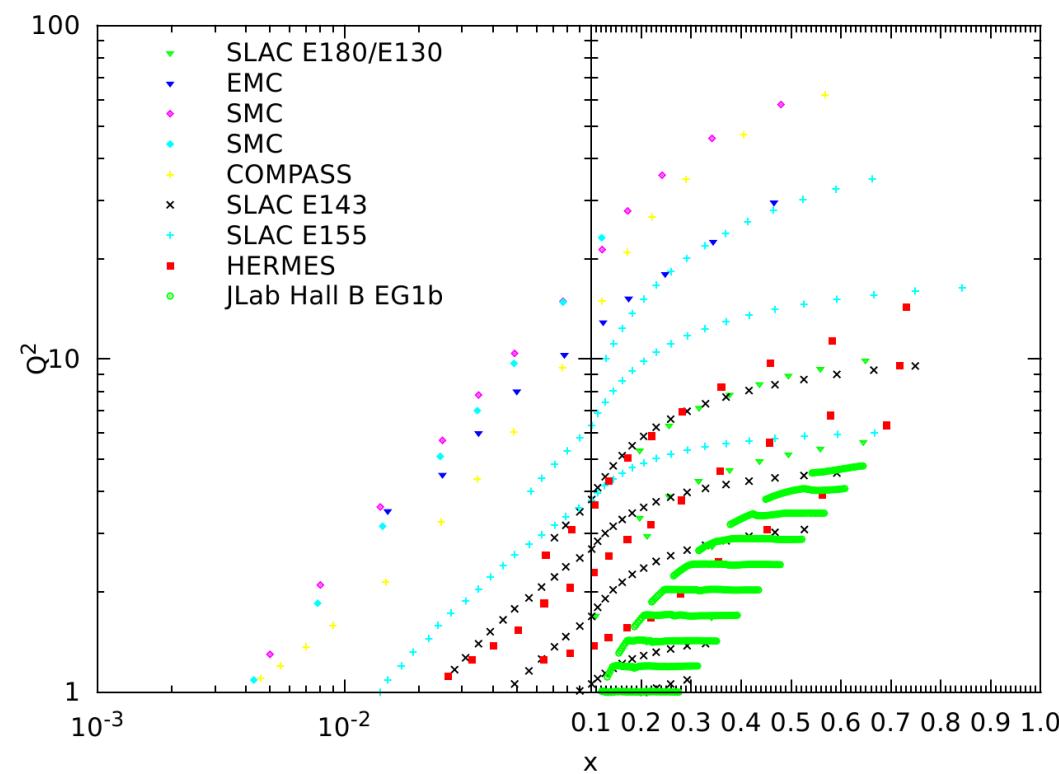
We consistently develop our own unpolarized analysis in parallel (JR)

Dedicated analyses of (some) JLab data: HALLAk, HALLAs, HALLB

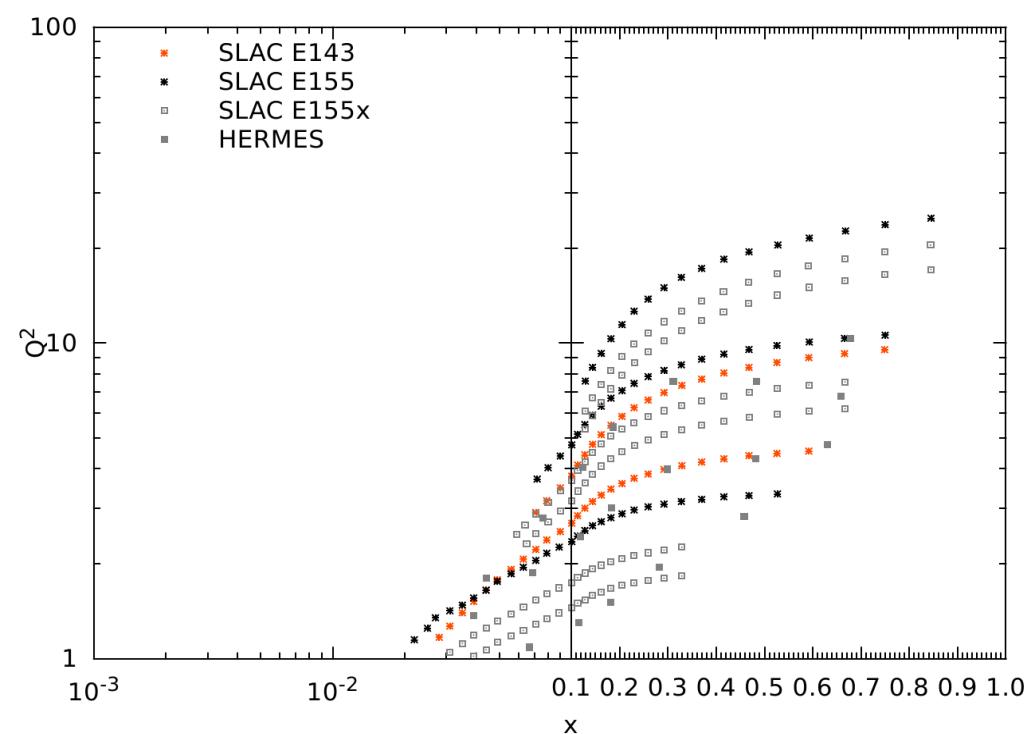
Data considered at this (first) stage

Minimal kinematic cuts imposed: $Q^2 \geq 1 \text{ GeV}^2$ $W^2 \geq 3.5 \text{ GeV}^2$

$A_{||}^p$ or A_1^p



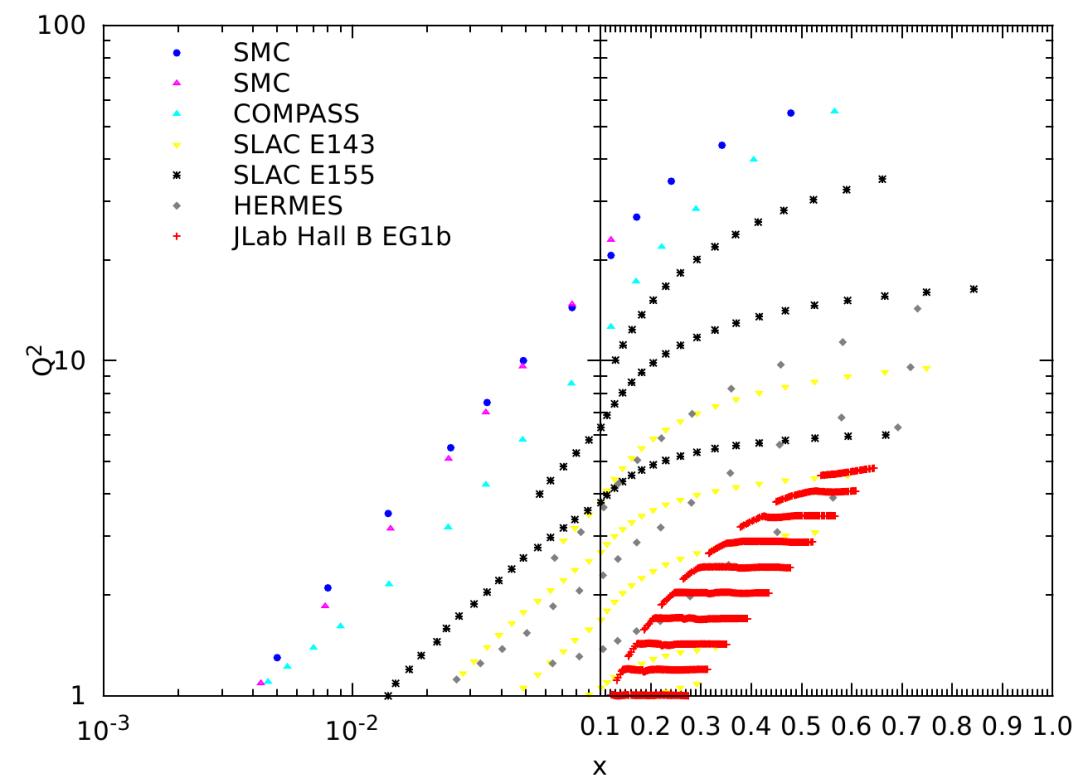
A_{\perp}^p or A_2^p



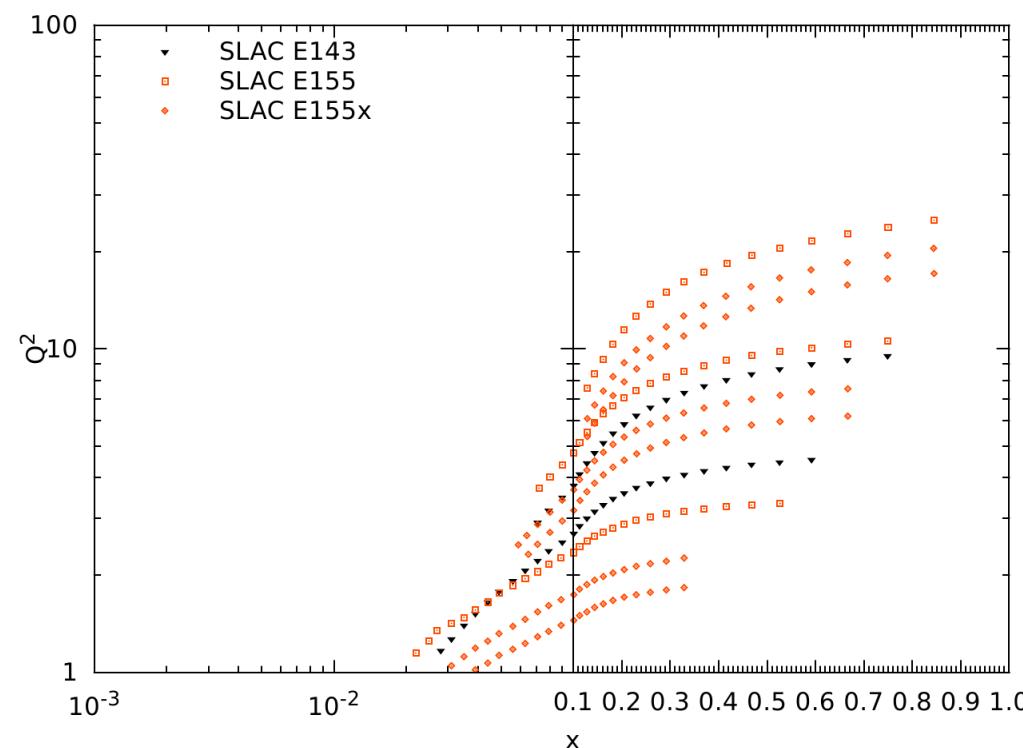
Data considered at this (first) stage

Minimal kinematic cuts imposed: $Q^2 \geq 1 \text{ GeV}^2$ $W^2 \geq 3.5 \text{ GeV}^2$

A_{\parallel}^d or A_1^d



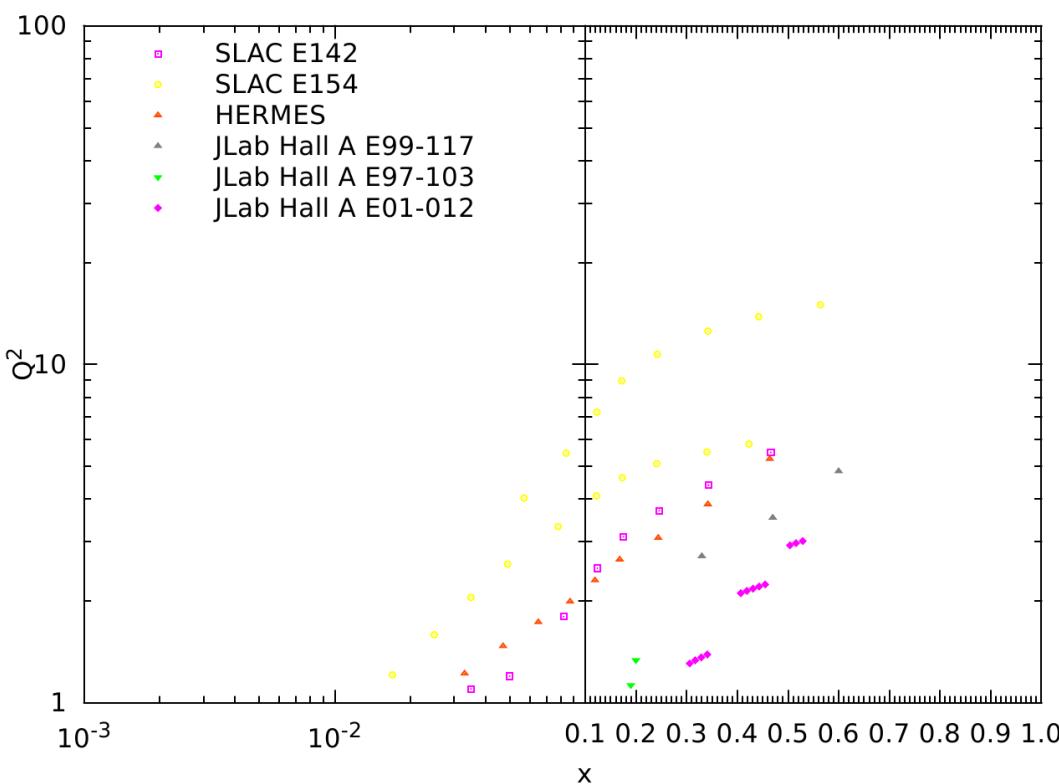
A_{\perp}^d or A_2^d



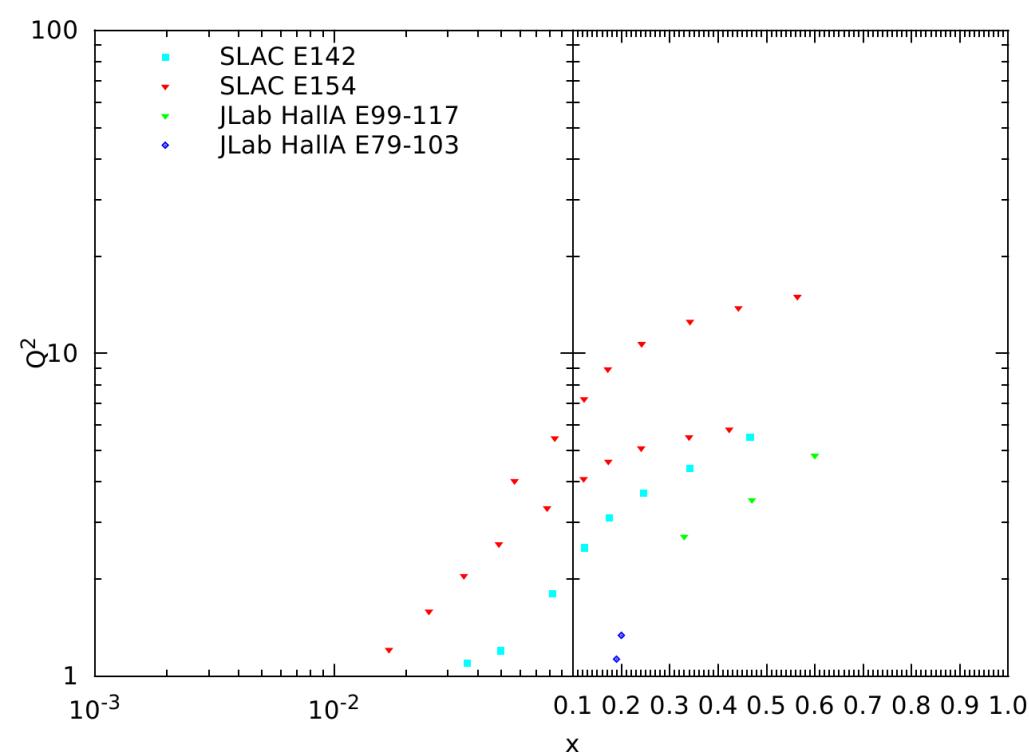
Data considered at this (first) stage

Minimal kinematic cuts imposed: $Q^2 \geq 1 \text{ GeV}^2$ $W^2 \geq 3.5 \text{ GeV}^2$

$A_{\parallel}^{3\text{He}}$ or $A_1^{3\text{He}}$



$A_{\perp}^{3\text{He}}$ or $A_2^{3\text{He}}$



Underlying QCD description

$$A_1 = (g_1 - \gamma^2 g_2) \frac{2x}{(1 + \gamma^2)F_2 - F_L} \quad A_2 = \gamma(g_1 + g_2) \frac{2x}{(1 + \gamma^2)F_2 - F_L}$$

NLO QCD calculation of (twist two) structure functions in Mellin space:

$$g_1(n, Q^2) = \frac{1}{2} \sum_{q, \bar{q}} e_q^2 (\Delta C_{qq}^1 \Delta q + \Delta C_g^1 \Delta g)$$

$$g_2(n, Q^2) = g_2^{WW} = -\frac{n-1}{n} g_1(n, Q^2)$$

RGE evolution of helicity parton distributions implemented as well in the space of complex moments (truncated solutions)

$$f(n) = \int_0^1 dx \ x^{n-1} f(x)$$

Parametrization

Only *two* independent combinations of quark distributions contribute:

$$x\Delta u^+(x, \mu_0^2) = N_u x^{a_u} (1-x)^{b_u} (1 + B_u x)$$

$$x\Delta d^+(x, \mu_0^2) = N_d x^{a_d} (1-x)^{b_d} (1 + B_d x)$$

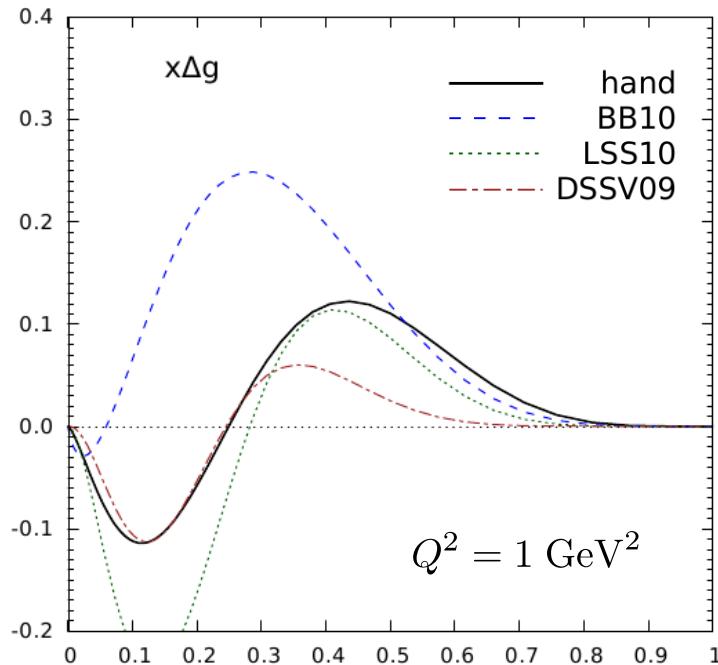
$$\Delta q^+ \equiv \Delta q + \Delta \bar{q}$$

Constraints from hyperon decays relate N_u and N_d , and fix N_s :

$$\int_0^1 (\Delta u^+ - \Delta d^+) dx = 1.269 \pm 0.003$$

$$\int_0^1 (\Delta u^+ + \Delta d^+ - 2\Delta s^+) dx = 0.586 \pm 0.031$$

The x -dependence of the sea has been fixed by hand (e.g. counting rules)



Formally Δg enters at second order,
but in practice our fits are *insensitive*:

$$\Delta\chi^2 \ll \Delta\chi^2_{1\sigma}$$

It has been fixed to a reasonable function (hand)

Statistical estimation

Least-squares estimator with a complete treatment of the systematic uncertainties (equivalent to the standard correlation matrix approach) [CTEQ]:

$$\chi^2 = \sum_{i=1}^N \frac{1}{\Delta_i^2} \left(D_i + \sum_{j=1}^M r_j \Delta_{ji} - T_i \right)^2 + \sum_{j=1}^M r_j^2$$

Here the systematic shifts are calculated *analytically*

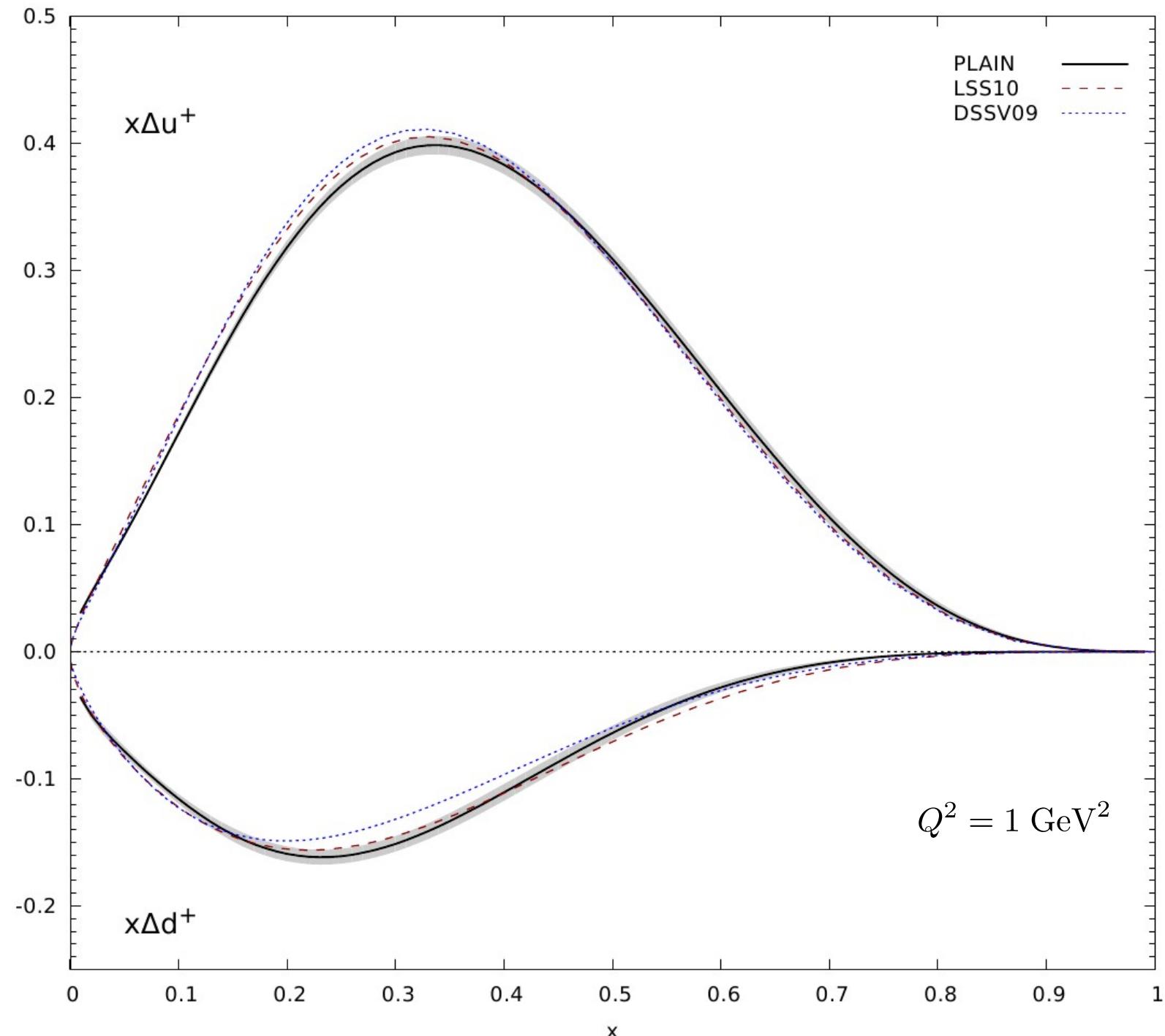
Unfortunately many experiments do not provide enough information

Errors estimated with the *Hessian* approach (linear propagation):

“Vicinity” of the minimum (tolerance) characterized by:

$$\Delta\chi^2 = \chi^2(p) - \chi^2(p^0) \leq T^2 = 1$$

Simple fit without further corrections: PLAIN



Improved description of nuclear targets: NSMEAR

In PLAIN the “effective polarizations” approximation have been used:

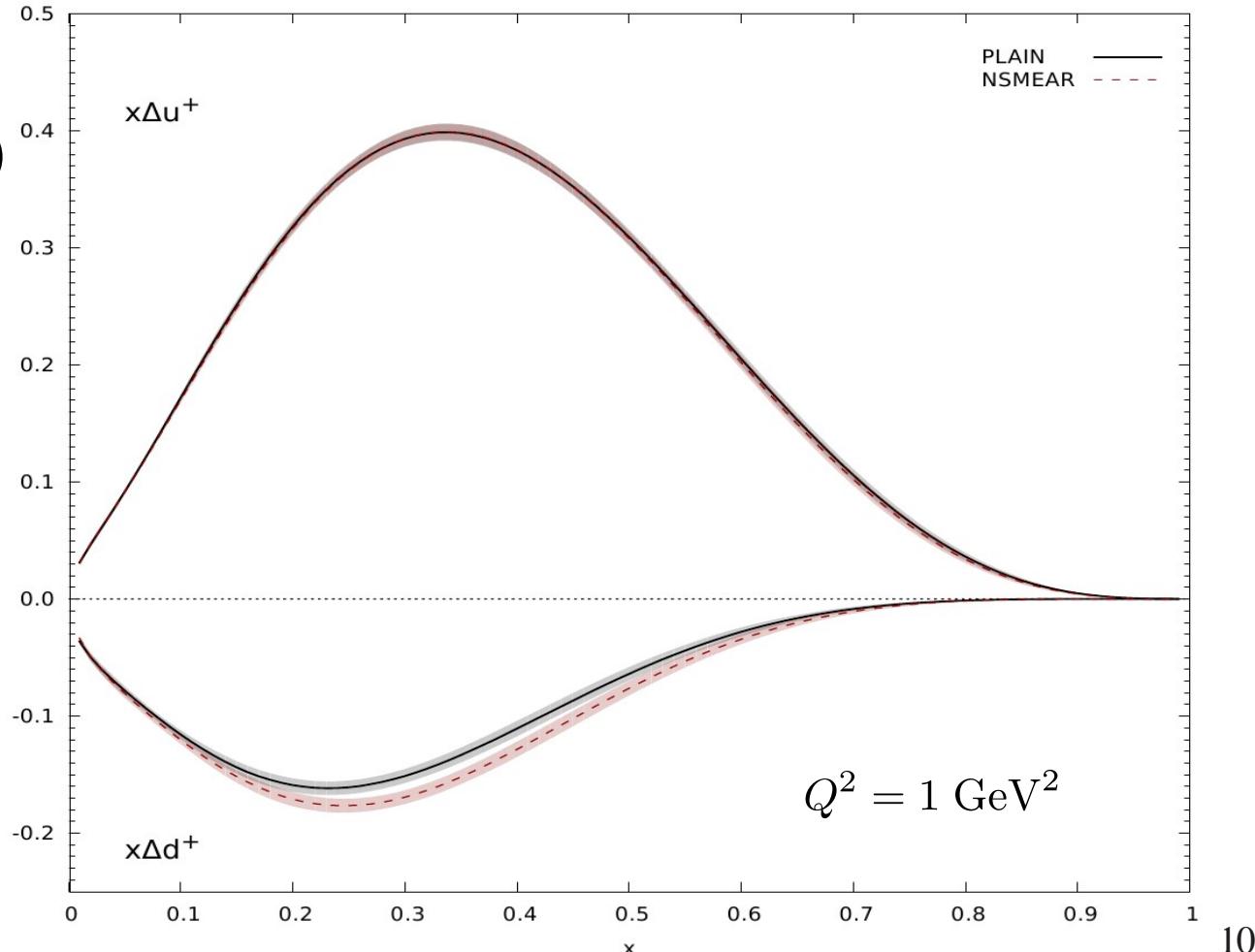
$$g_1^d = (1 - \frac{3}{2}0.06)(g_1^p + g_1^n) \quad g_1^{He3} = 0.86 g_1^n - 0.059 g_1^p$$

An improvement over this is to consider “smearing” functions derived from the nuclear wave functions

$$g_i^A(x) = \sum_{\substack{j=1,2 \\ N=p,n}} \int dy f_{jN}(y, \gamma) g_j^N(\frac{x}{y})$$
$$\gamma^2 = 1 + 4 \frac{M^2}{Q^2} x^2$$

Relevant for Δd in the medium- to large- x region

(This approach is also used in the unpolarized part)



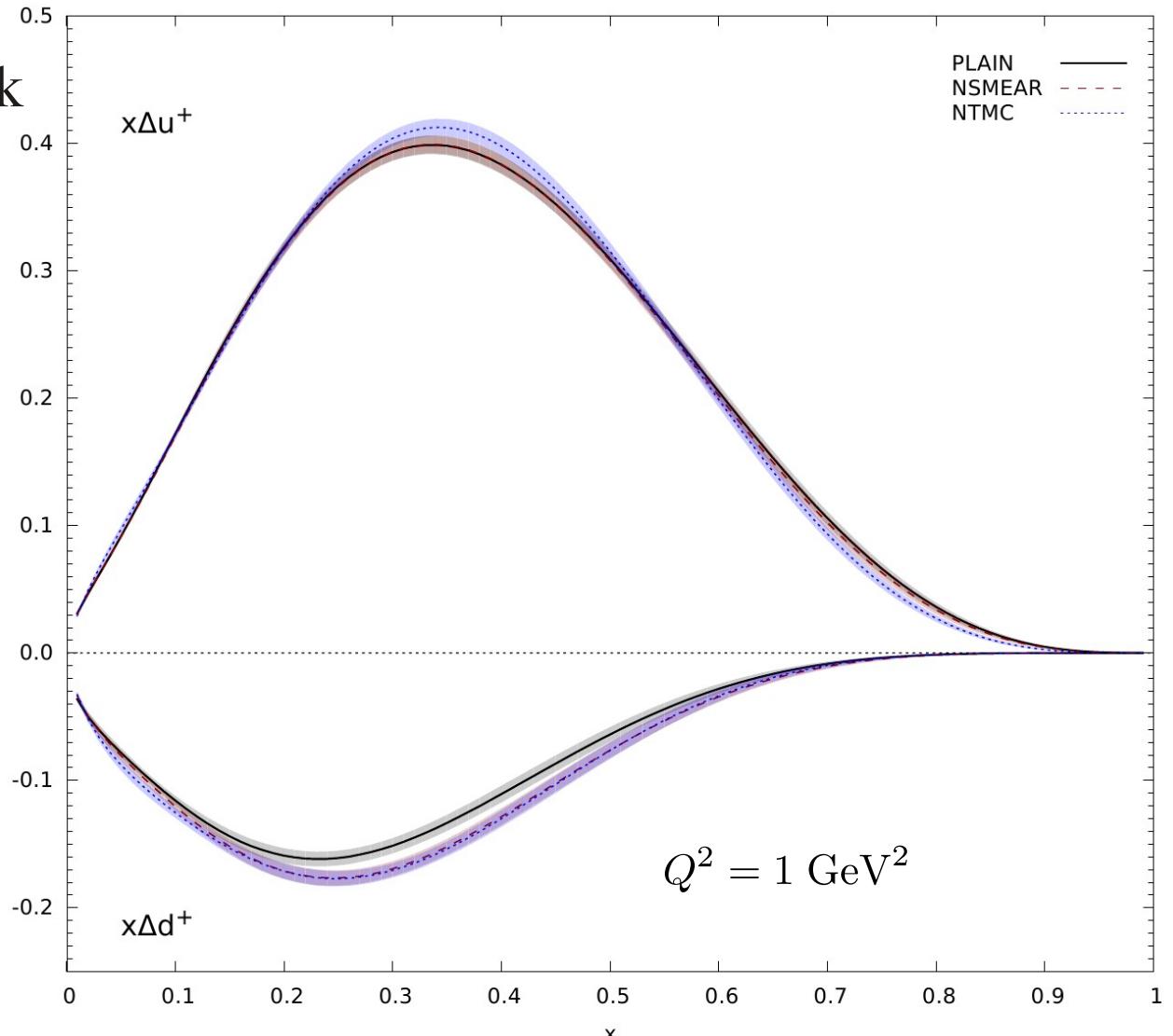
Plus target-mass corrections: NTMC

We use power corrections from finite target mass calculated in the OPE approach:

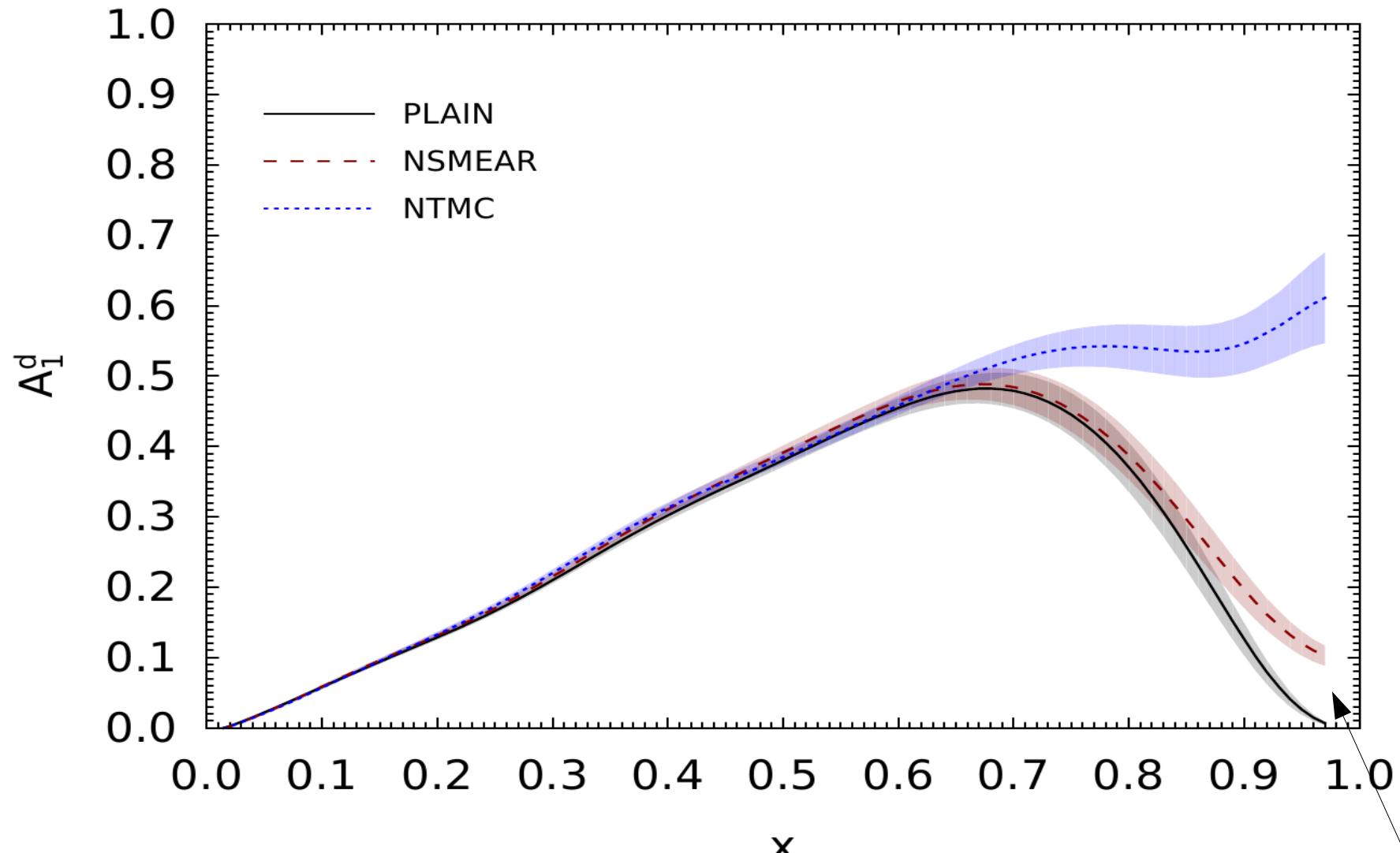
$$g_1^{\text{TMC}}(n) = g_1(n) + \frac{M^2}{Q^2} \frac{n^2(n+1)}{(n+2)^2} g_1(n+2) + \frac{M^4}{Q^4} \dots + \mathcal{O}\left(\frac{M^8}{Q^8}\right)$$

where the Wandzura-Wilczek relation holds also after
TMCs [Bluemlein, Tkabladze 99]

These are relevant for Δu
in the medium- to
large- x region



Plus target-mass corrections: NTMC



Clearly they are needed for the asymmetries (TMCs included in the unpolarized SFs!):

Higher twist contributions

We consider also corrections from higher twist contributions:

$$g_1 = g_1^{\tau=2} + g_1^{\tau=3} + g_1^{\tau=4}$$

$$g_2 = g_2^{\tau=2} + g_2^{\tau=3}$$

The Bluemlein-Tkabladze relation: $g_1^{\tau=3}(x, Q^2) = 4x^2 \frac{M^2}{Q^2} \left(g_2^{\tau=3}(x, Q^2) - 2 \int_x^1 \frac{dy}{y} g_2^{\tau=3}(y, Q^2) \right)$

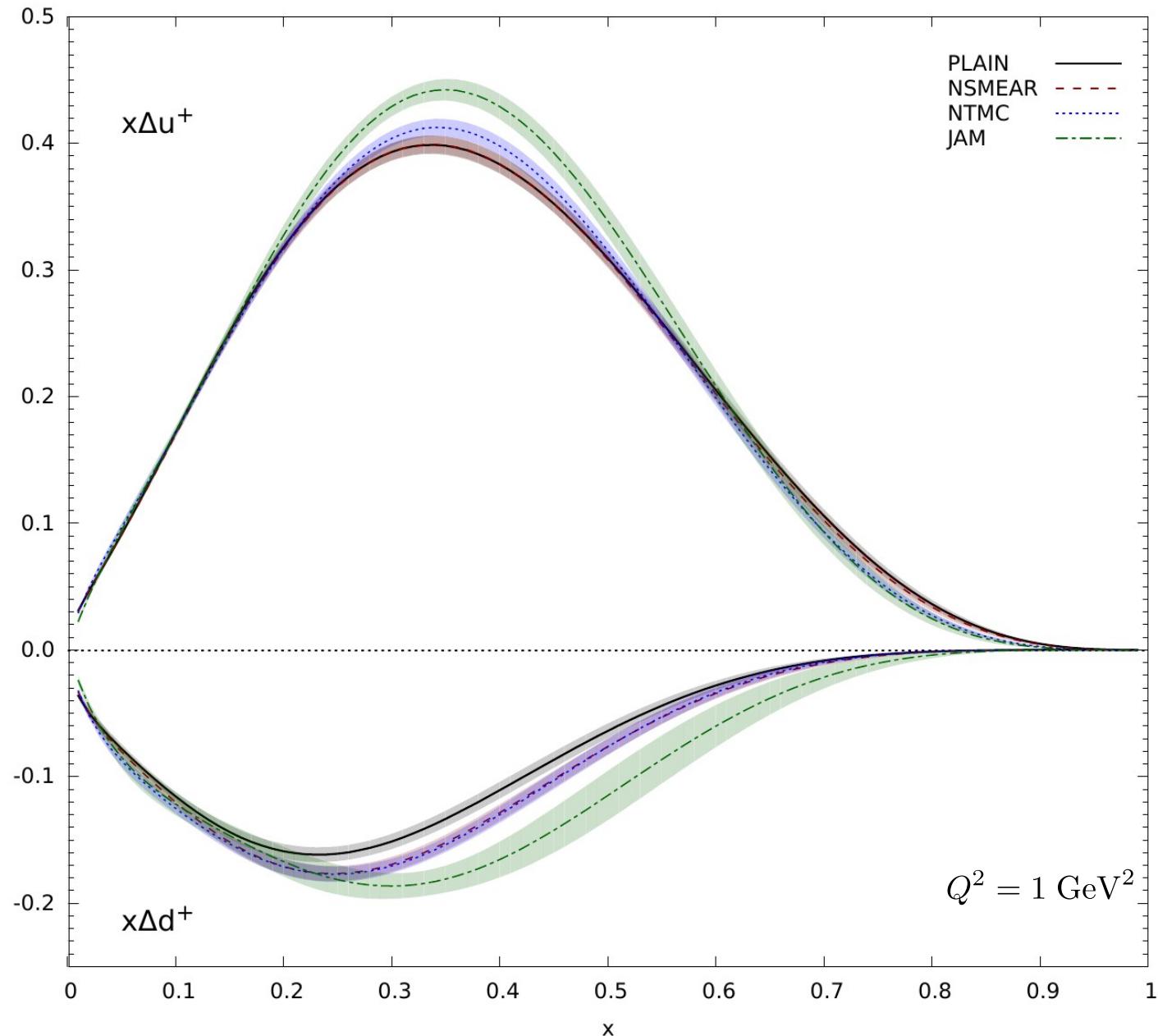
With a phenomenological parametrization [Braun *et al.* 09]:

$$g_2^{\tau=3} = A[\ln x + (1-x) + \frac{1}{2}(1-x)^2] + (1-x)^3[B + C(1-x) + D(1-x)^2 + E(1-x)^3]$$

And a splines approximation for: $g_1^{\tau=4} = \frac{h(x)}{Q^2}$ $\int_0^1 dx g_2^{\tau=3} = 0$

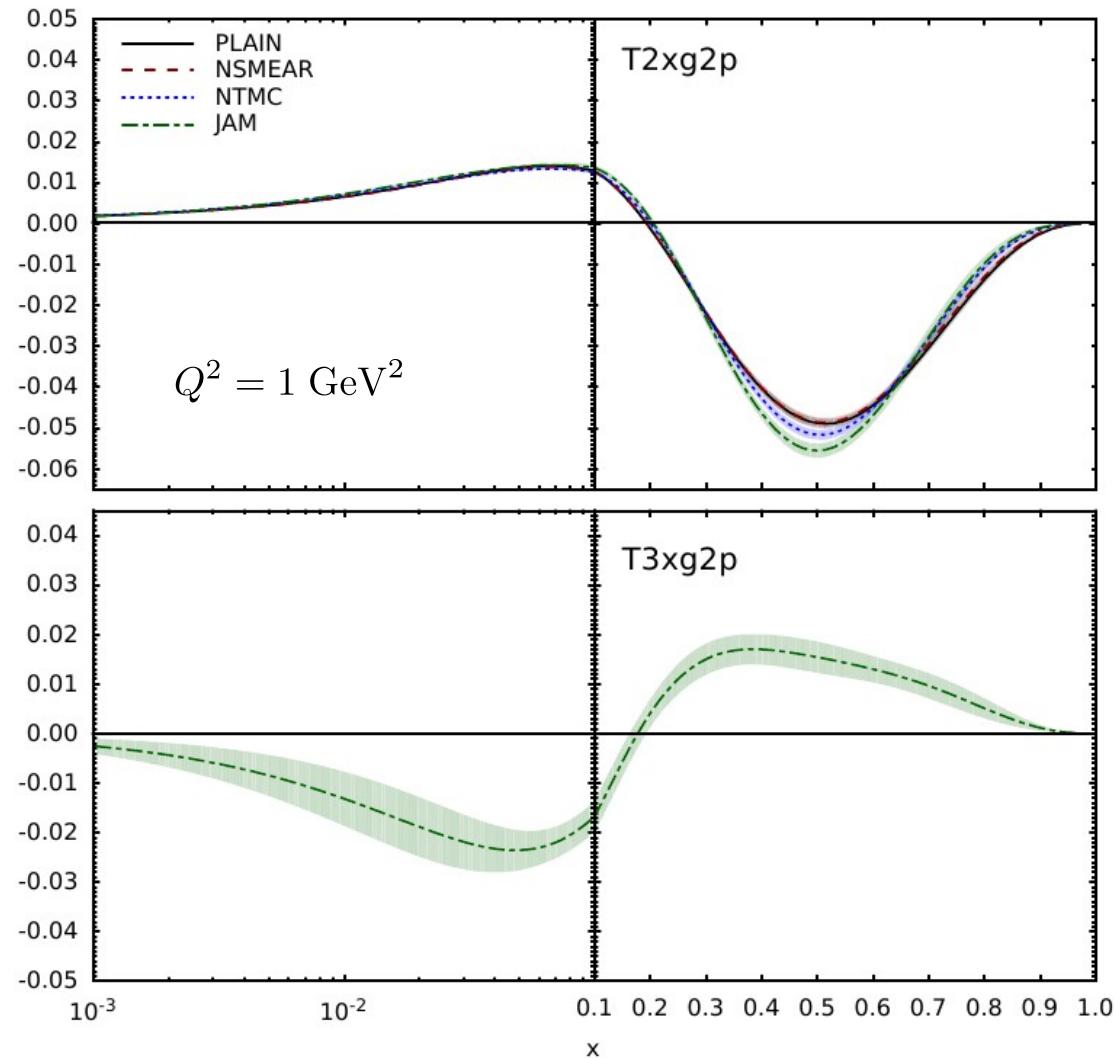
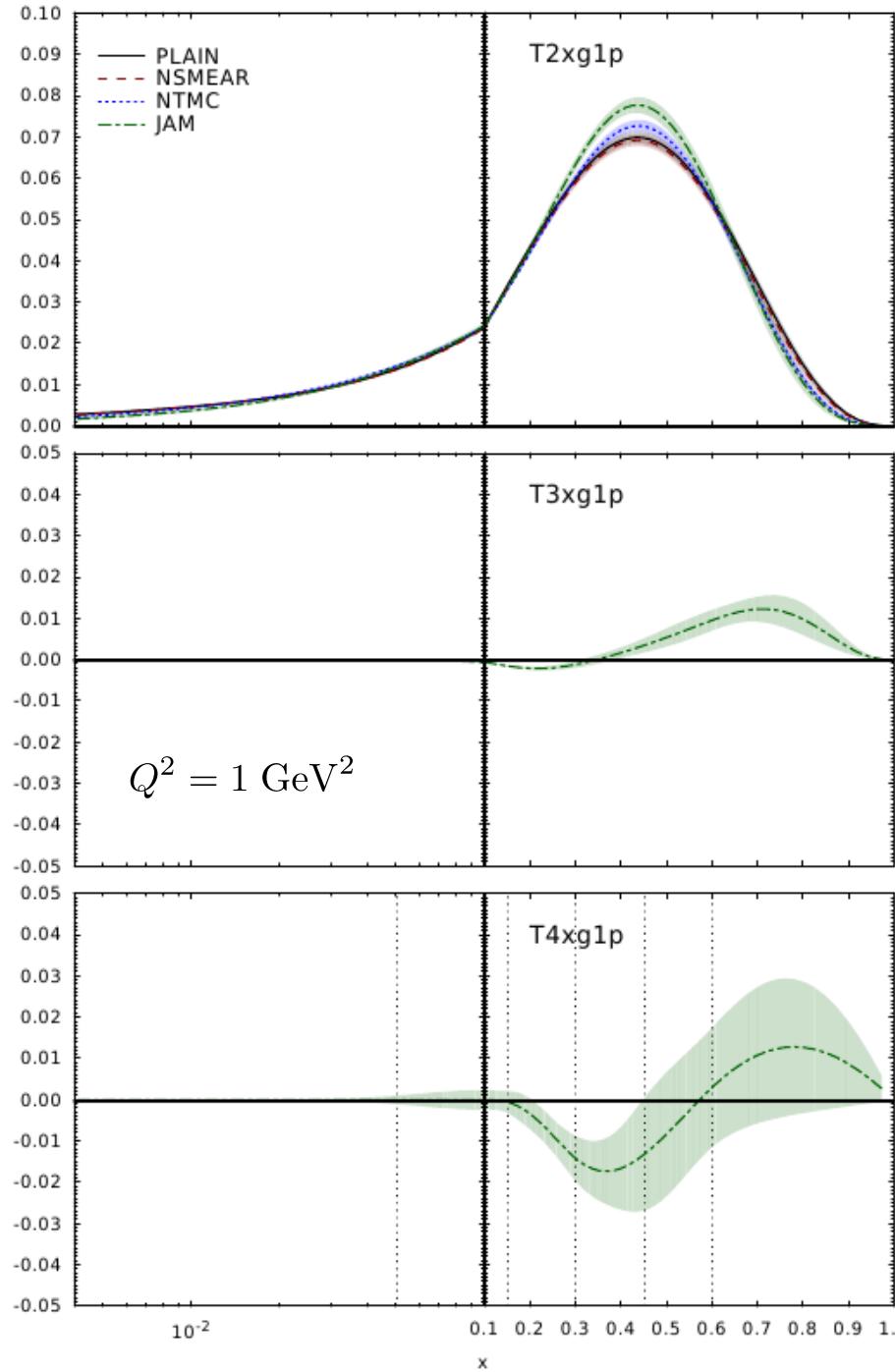
Possible scale dependence in h and $g_2^{\tau=3}$ have been neglected

Including all these corrections: JAM

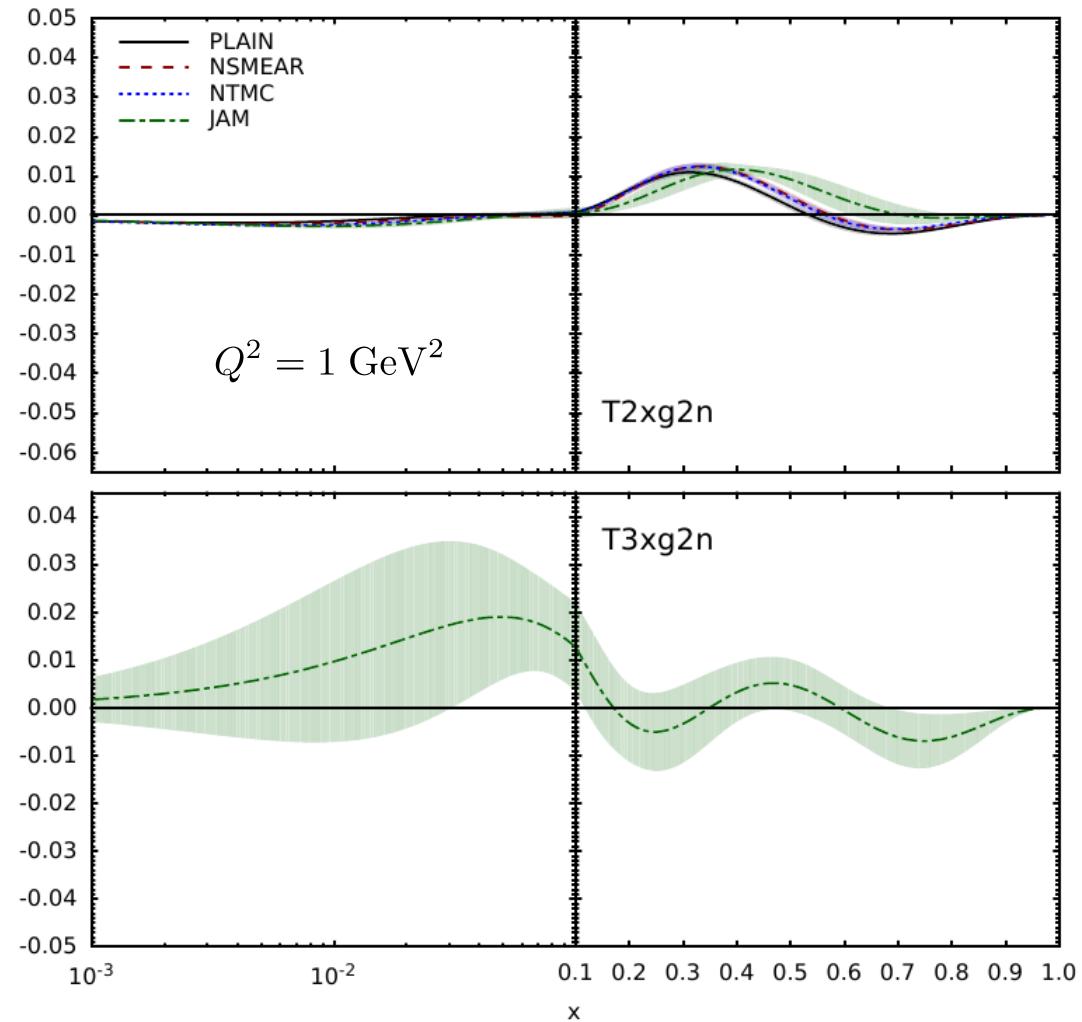
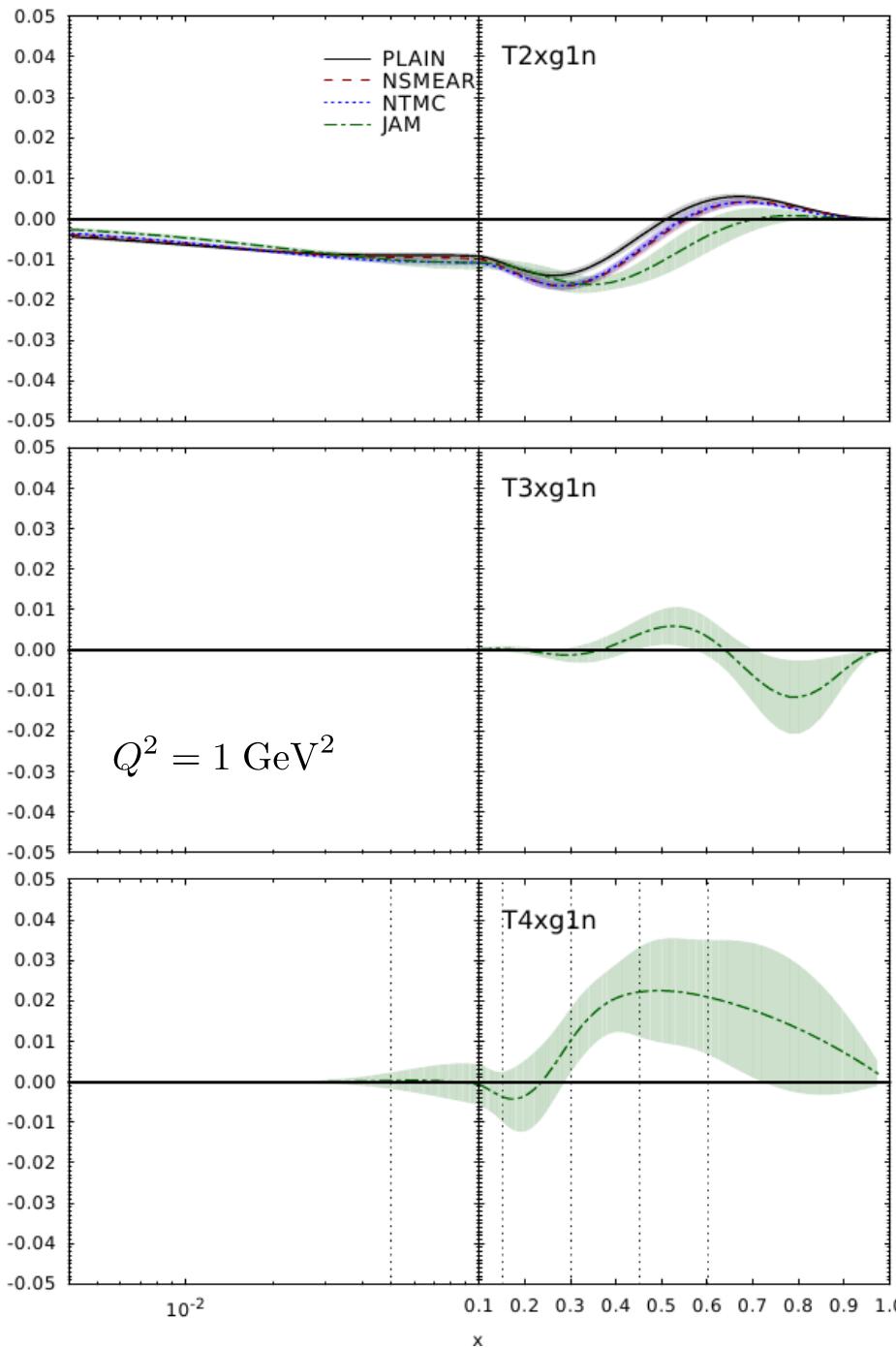


Manifest importance of higher twist contributions

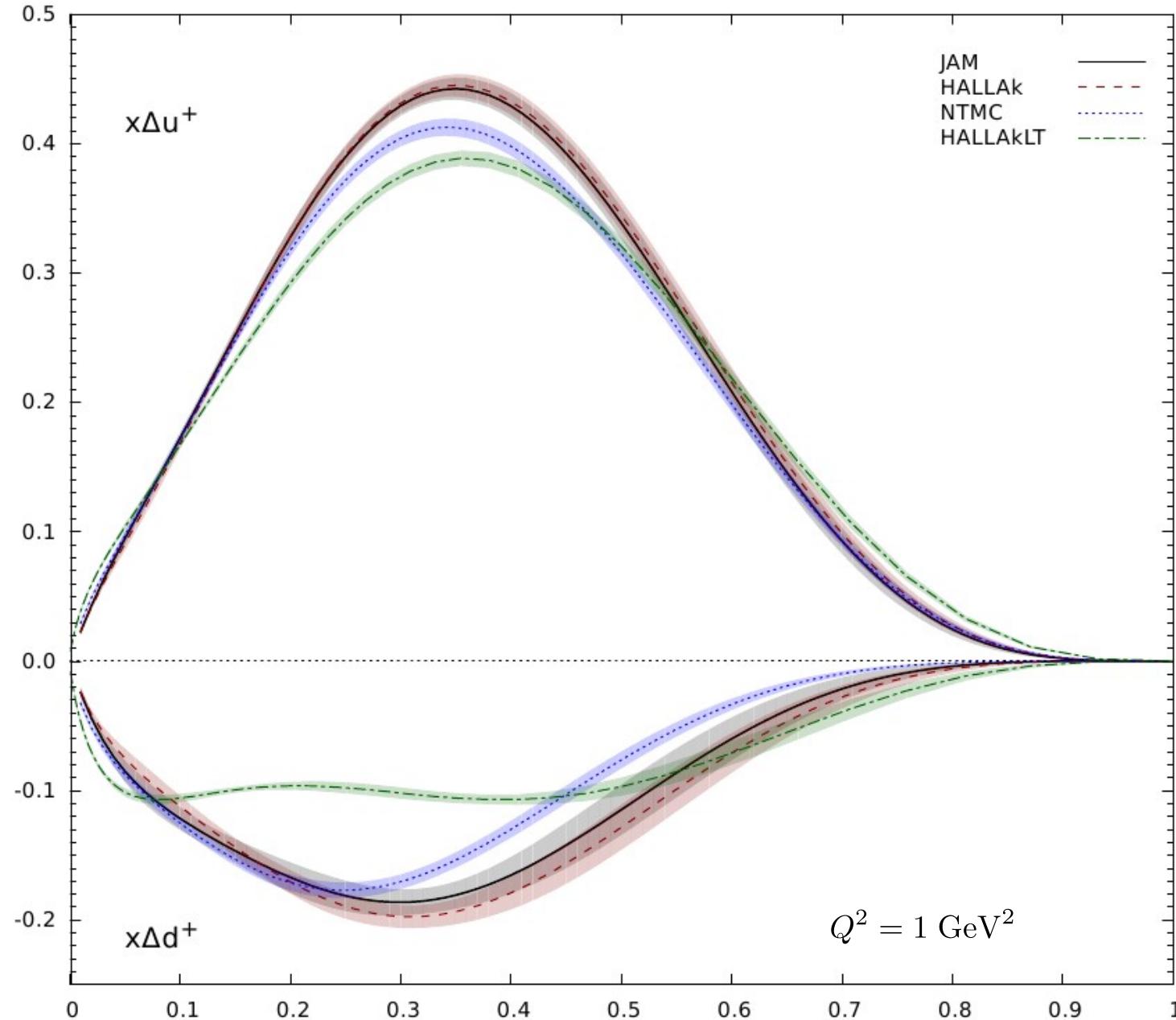
Including all these corrections: JAM



Including all these corrections: JAM

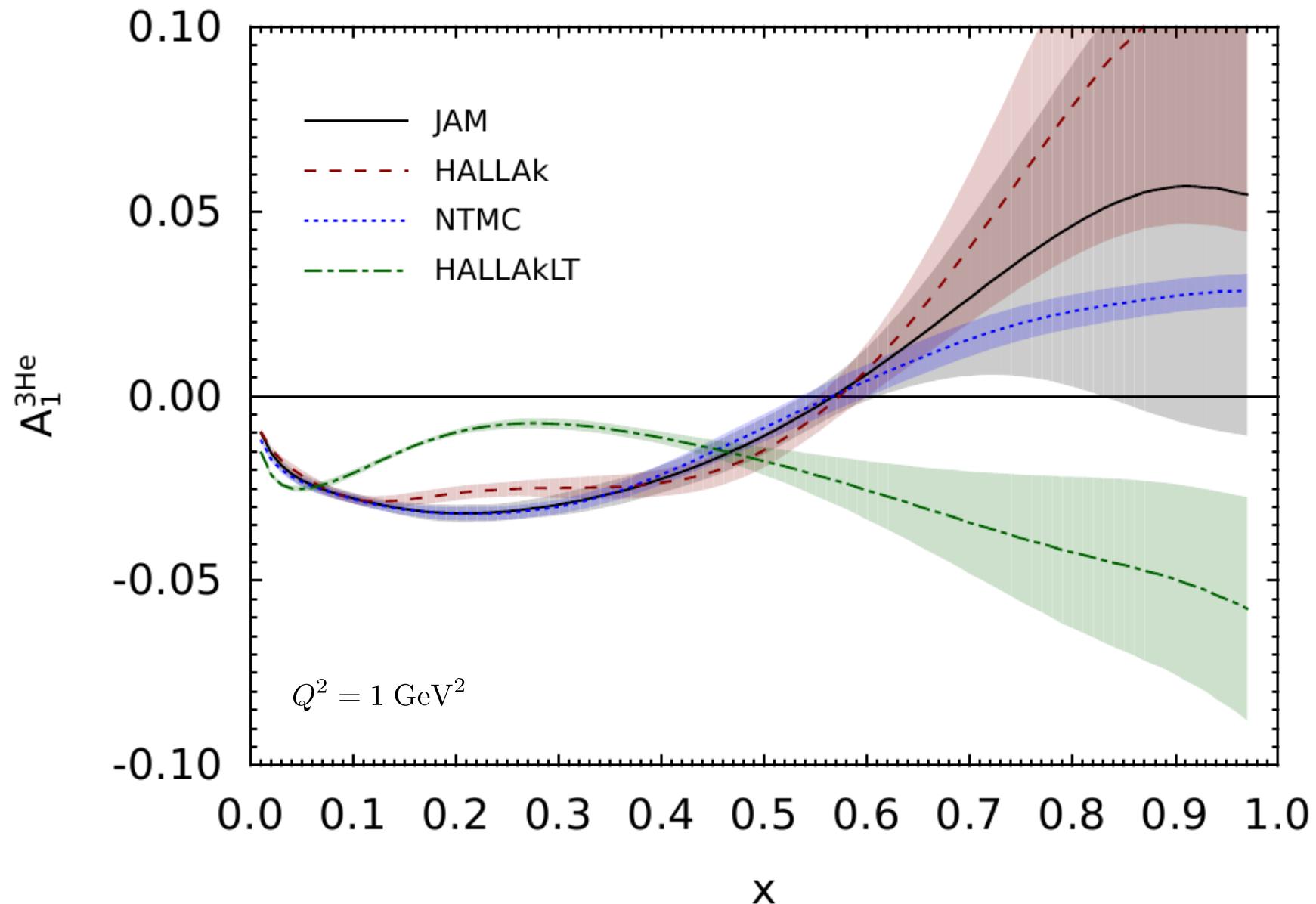


Including Jlab Hall A E97-103: HALLAk

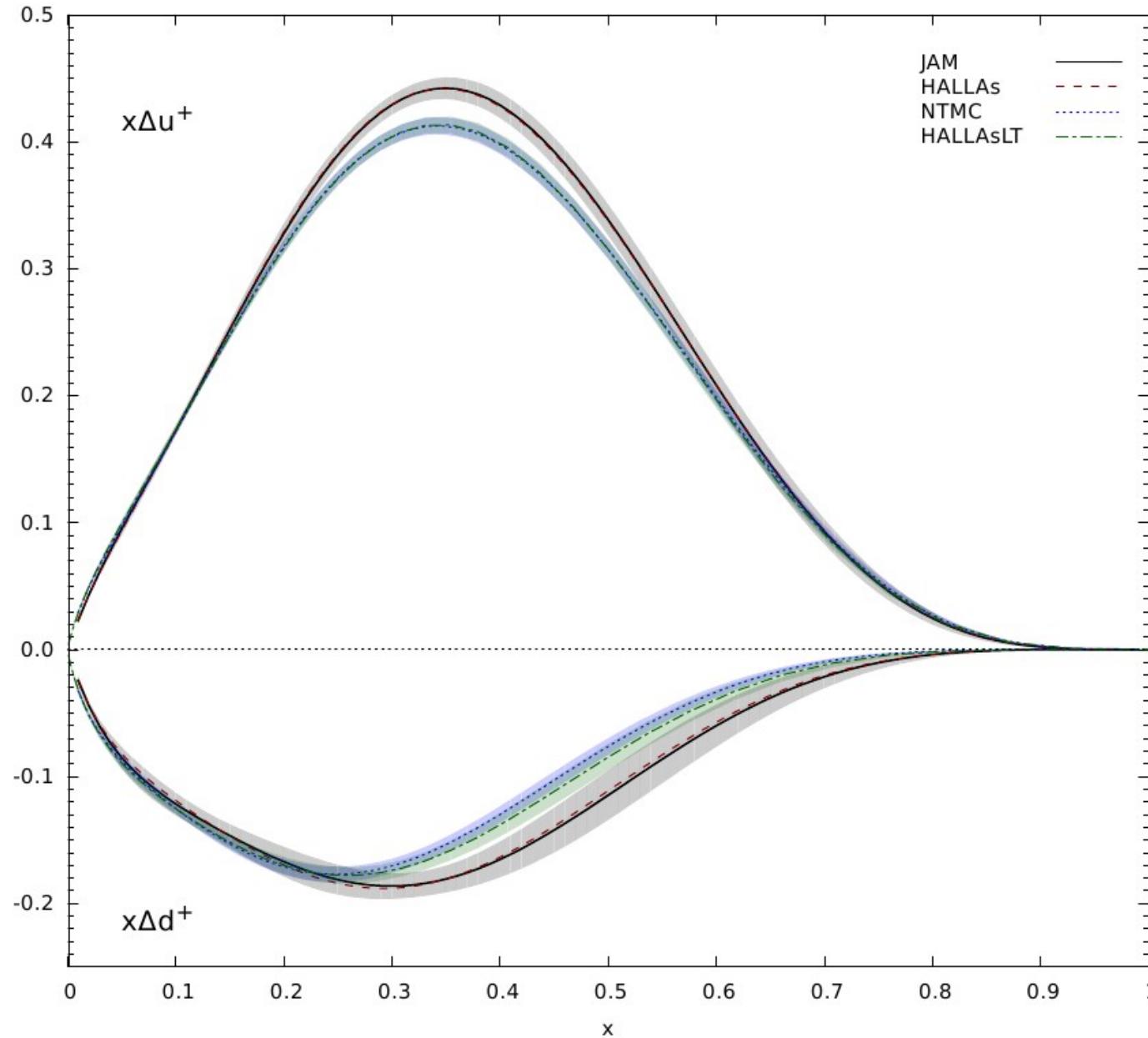


Better with HT, but χ^2 very large (for $A_{||} \stackrel{x}{\text{LT}}$ about 15, HT about 4)

Including Jlab Hall A E97-103: HALLAk



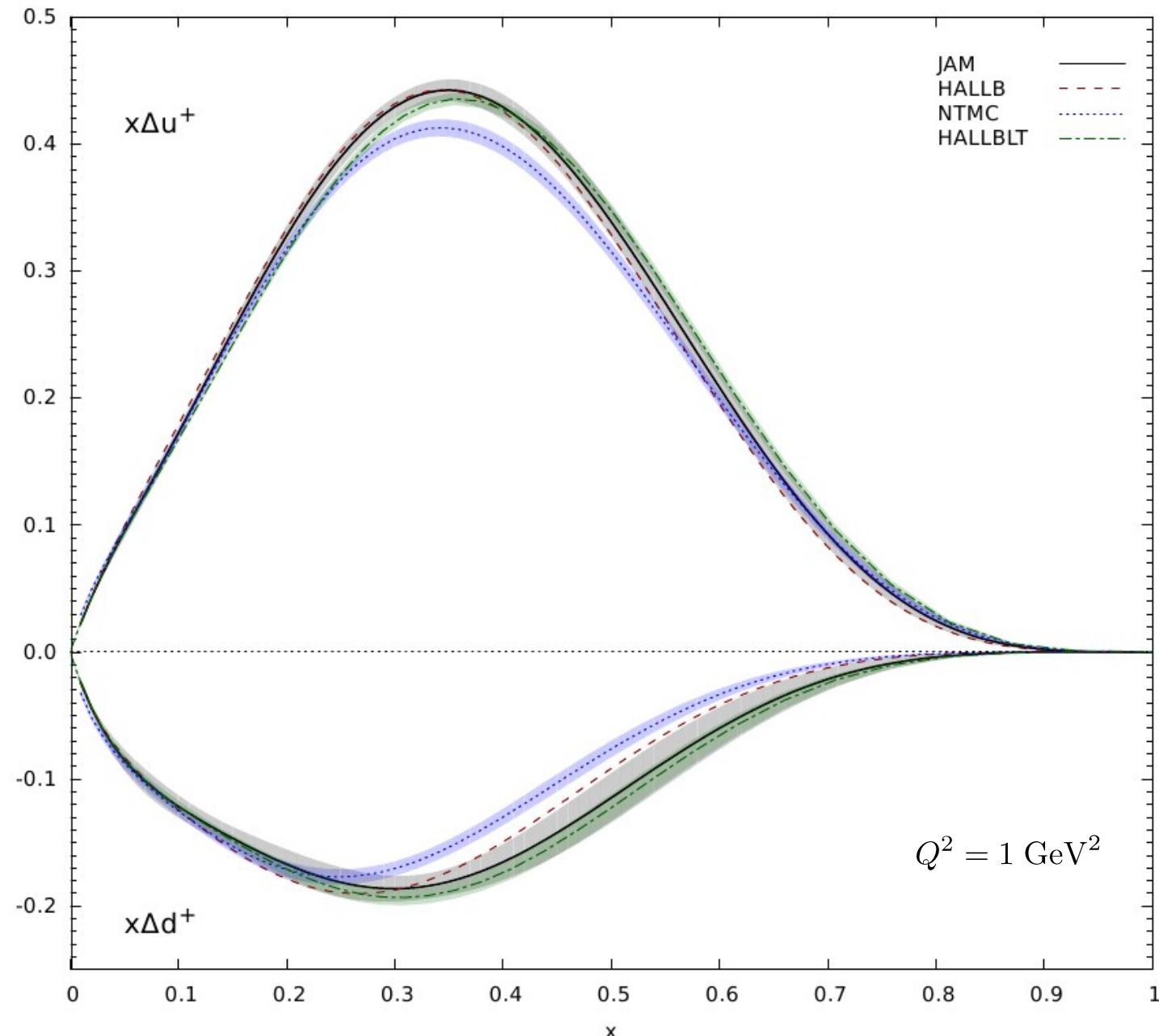
Including Jlab Hall A E01-012: HALLAs



Central values do not move that much, but χ^2 rather large (LT about 3, HT about 2.5)

Including Jlab Hall B: HALLB

Are they ready? A combination of the data points would be great



Summary and outlook

First JAM results practically *ready*

Some things still to be done: $x \rightarrow 1$ studies, OAM, symmetric sea, ...

We will present them in the meetings in Denver (GHP + APS) ,
and in DIS 2013 (Alberto)

It is time to write our first paper!

Some tasks to which you could contribute?

- Help preparing tables and figures, e.g. comparisons with A_{\parallel} and A_{\perp} , combinations of data, and such (I can provide you with our predictions)
- Help with the text? Especially in what refers to your own data (if it applies)
- Suggestions and comments are also welcome